

Dynamics near a nonhyperbolic critical pt in \mathbb{R}^2

(M, Chapter 6.1–6.3)

$$\frac{dX}{dt} = AX + g(X), \quad A^{\text{def}}, |g(x)| \leq C|x|^2$$

A — nonhyperbolic:

① $\lambda_1 = 0, \lambda_2 \neq 0$

② $\lambda_1 = \lambda_2 = 0$, but only one eigenvector

$$A = \begin{bmatrix} 0 & 0 \\ 0 & \lambda_2 \end{bmatrix}$$

$$A = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

③ $\lambda_1 = \lambda_2 = 0$, two eigenvectors ④ $\lambda_1 = iw, \lambda_2 = -iw$

$$A = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix}$$

$$A = \begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix}$$

Case ③ $A = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix}$

$$\left\{ \begin{array}{l} \frac{dx}{dt} = P(x, y), \quad |P(x, y), Q(x, y)| \leq C(|x|^2 + |y|^2) \\ \frac{dy}{dx} = Q(x, y), \end{array} \right.$$

↓ Taylor-expansion

$$P(x, y) = \sum_{i=0}^2 a_{2i} x^i y^{2-i} + \sum_{i=0}^3 a_{3i} x^i y^{3-i} + \dots + \sum_{i=0}^n a_{ni} x^i y^{n-i} + \dots$$

$$Q(x, y) = \underbrace{\sum_{i=0}^2 b_{2i} x^i y^{2-i}}_{\text{homogeneous degree } = 2} + \underbrace{\sum_{i=0}^3 b_{3i} x^i y^{3-i}}_3 + \dots + \underbrace{\sum_{i=0}^n b_{ni} x^i y^{n-i}}_n + \dots$$

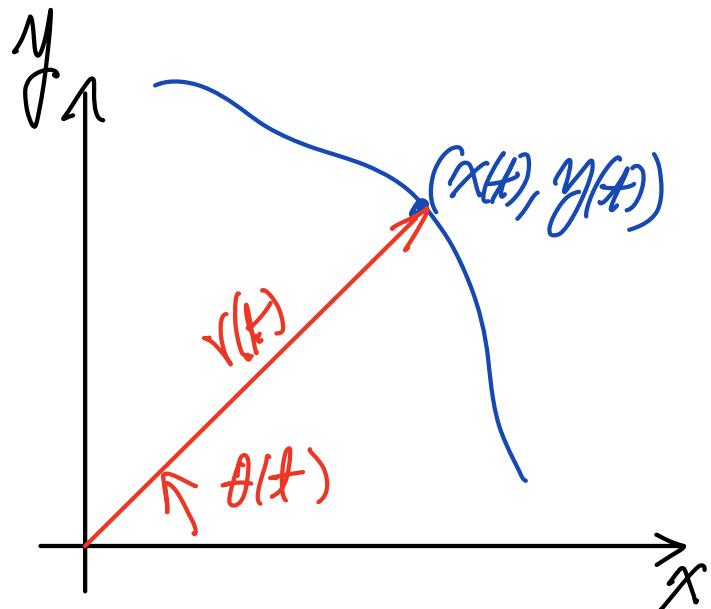
homogeneous degree = 2

3

n

Use Polar Coordinates

$$\begin{cases} x = r \cos \theta \\ y = r \sin \theta, \end{cases} \quad \begin{cases} r^2 = x^2 + y^2 \\ \theta = \tan^{-1}\left(\frac{y}{x}\right) \end{cases}$$



$$\frac{d(r^2)}{dt} = 2x\dot{x} + 2y\dot{y} \Rightarrow \dot{r} = \frac{1}{r}(x\dot{x} + y\dot{y})$$

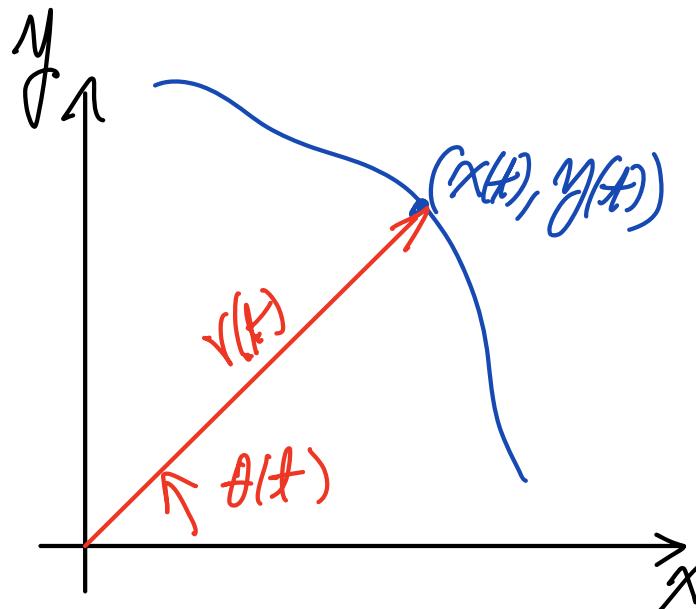
\approx
 $2\dot{r}r$

$$\frac{dr}{dt} = \frac{1}{r} [r \cos \theta P(r \cos \theta, r \sin \theta) + r \sin \theta Q(r \cos \theta, r \sin \theta)]$$

$$= (\cos \theta) P(r \cos \theta, r \sin \theta) + (\sin \theta) Q(r \cos \theta, r \sin \theta)$$

Use Polar Coordinates

$$\begin{cases} x = r \cos \theta \\ y = r \sin \theta, \end{cases} \quad \begin{cases} r^2 = x^2 + y^2 \\ \theta = \tan^{-1}\left(\frac{y}{x}\right) \end{cases}$$



$$\frac{d\theta}{dt} = \frac{d}{dt} \left[\tan^{-1}\left(\frac{y}{x}\right) \right] = \frac{1}{1 + \frac{y^2}{x^2}} \left(\frac{x\dot{y} - y\dot{x}}{x^2} \right) = \frac{x\dot{y} - y\dot{x}}{x^2 + y^2}$$

$$\frac{d\theta}{dt} = \frac{(r \cos \theta) Q(-) - (r \sin \theta) P(-)}{r^2}$$

$$= \frac{1}{r} \left[(\cos \theta) Q(r \cos \theta, r \sin \theta) - (\sin \theta) P(r \cos \theta, r \sin \theta) \right]$$

Equation in Polar Coordinate

$$\left\{ \begin{array}{l} \frac{dr}{dt} = (\cos\theta) P(r\cos\theta, r\sin\theta) + (\sin\theta) Q(r\cos\theta, r\sin\theta) \\ \frac{d\theta}{dt} = \frac{1}{r} [\cos\theta Q(r\cos\theta, r\sin\theta) - \sin\theta P(r\cos\theta, r\sin\theta)] \end{array} \right.$$

$$\frac{dr}{d\theta} = \frac{(\cos\theta) P(r\cos\theta, r\sin\theta) + (\sin\theta) Q(r\cos\theta, r\sin\theta)}{\frac{1}{r} [\cos\theta Q(r\cos\theta, r\sin\theta) - \sin\theta P(r\cos\theta, r\sin\theta)]}$$

Equation in Polar Coordinate

$$\left\{ \begin{array}{l} \frac{dr}{dt} = (\cos\theta) P(r\cos\theta, r\sin\theta) + (\sin\theta) Q(r\cos\theta, r\sin\theta) \\ \frac{d\theta}{dt} = \frac{1}{r} [\cos\theta Q(r\cos\theta, r\sin\theta) - \sin\theta P(r\cos\theta, r\sin\theta)] \end{array} \right.$$

Assume P, Q are homogeneous of degree n

$$P(\alpha x, \alpha y) = \underline{\alpha^n} P(x, y), \quad Q(\alpha x, \alpha y) = \underline{\alpha^n} Q(x, y)$$

Then $P(r\cos\theta, r\sin\theta) = r^n P(\cos\theta, \sin\theta), \dots$

Equation in Polar Coordinate

$$\frac{dr}{dt} = r^n [C P(c, s) + S Q(c, s)]$$

$$C = \cos \theta$$
$$S = \sin \theta$$

$$\frac{d\theta}{dt} = r^{n-1} [C Q(c, s) - S P(c, s)]$$

$$\frac{dr}{d\theta} = r \left[\frac{C P(c, s) + S Q(c, s)}{S Q(c, s) - C P(c, s)} \right]$$

$$g(\theta)$$

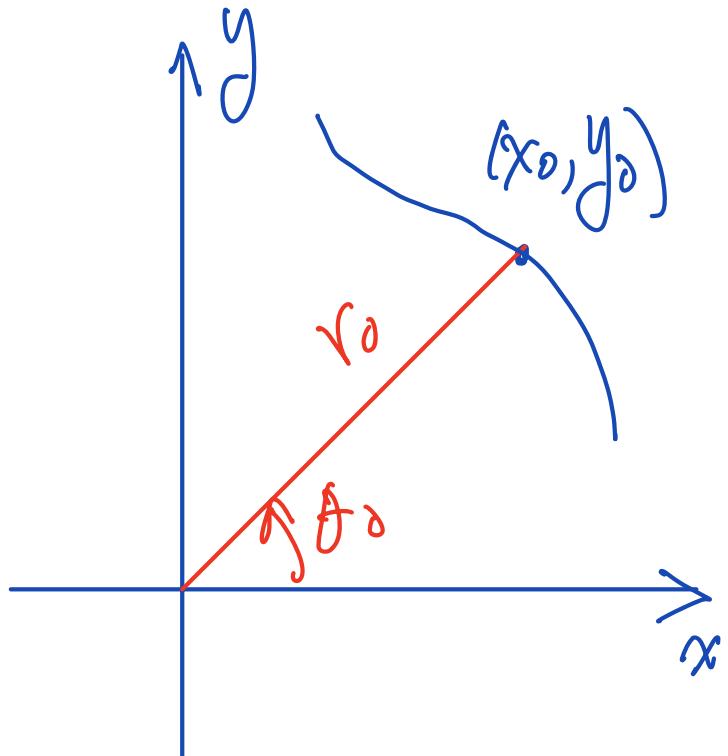
Solution in Polar Coordinate

- $\frac{dr}{d\theta} = r g(\theta) \leftarrow \text{separable equation}$

- $\int_{r_0}^r \frac{dr}{r} = \int_{\theta_0}^{\theta} g(\theta) d\theta$

- $\ln r = \ln r_0 + \int_{\theta_0}^{\theta} g(\varphi) d\varphi$

- $r(\theta) = r(\theta_0) e^{\int_{\theta_0}^{\theta} g(\varphi) d\varphi}$ (e.g. $\theta_0 = 0$)



Solution in Polar Coordinate

- $\frac{dr}{d\theta} = r g(\theta) \leftarrow \text{separable equation}$

- $\int_{r_0}^r \frac{dr}{r} = \int_{\theta_0}^{\theta} g(\theta) d\theta$ Note:
 $g(\varphi) = \varphi (\varphi + 2\pi)$

- $\ln r = \ln r_0 + \int_{\theta_0}^{\theta} g(\varphi) d\varphi$ Define:

$$G_I = \int_0^{2\pi} g(\varphi) d\varphi$$

- $r(\theta) = r(\theta_0) e^{\int_{\theta_0}^{\theta} g(\varphi) d\varphi}$

Solution in Polar Coordinate

$$r(\theta) = r(\theta_0) e^{\int_{\theta_0}^{\theta} g(\varphi) d\varphi}$$

Note :

$$g(\varphi) = \varphi (\varphi + 2\pi)$$

assume well defined

Define:

$$r(\theta_0 + 2\pi) = r(\theta_0) e^{G_1}$$

$$r(\theta_0 + 2\pi/2) = r(\theta_0) e^{2G_1}$$

$$r(\theta_0 + 2\pi/3) = r(\theta_0) e^{3G_1}$$

⋮ ⋮

$$r(\theta_0 + 2\pi n) = r(\theta_0) e^{nG_1}$$

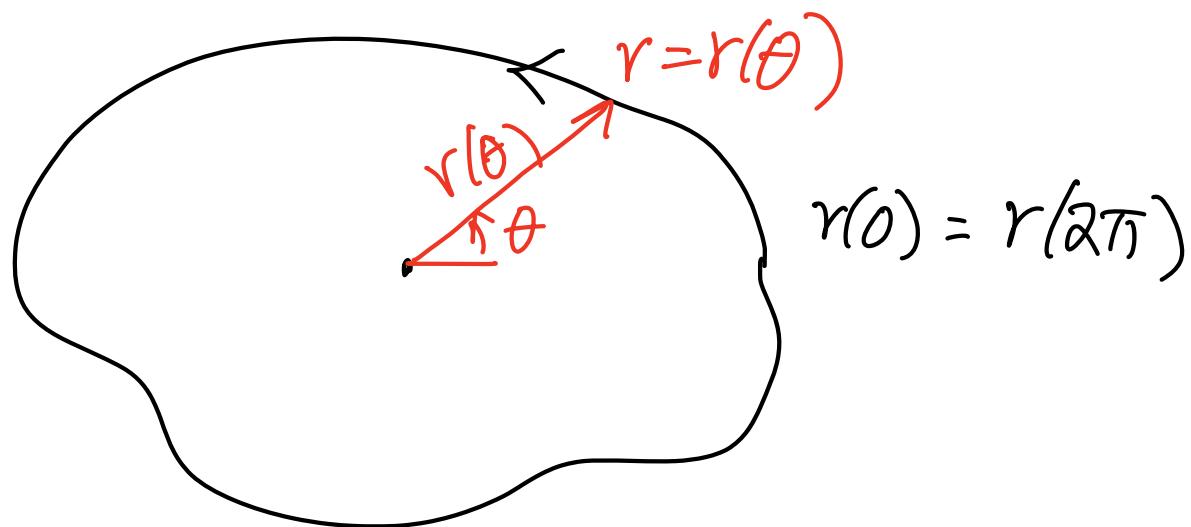
$\xrightarrow{n \rightarrow \infty} ?$

$$G_1 = \int_0^{2\pi} g(\varphi) d\varphi$$

Solution in Polar Coordinate

Suppose $G_1 = \int_0^{2\pi} g(\varphi) d\varphi$ is well-defined.

If $G_1 = 0$, then $r(\theta=0) = r(\theta=2\pi)$
i.e. periodic orbit.



Solution in Polar Coordinate

Suppose $G = \int_0^{2\pi} g(\varphi) d\varphi$ is well-defined.

If $G < 0$, then $\ln r(2\pi) = \ln r(0) + G$

$$r(2\pi) = r(0) e^G$$

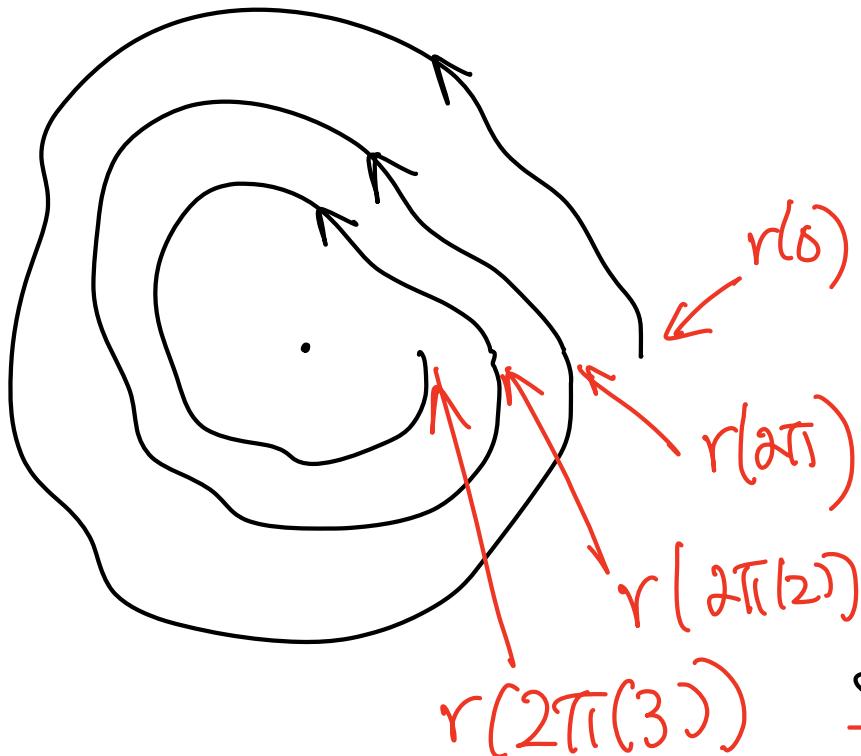
$$r(2\pi_2) = r(0) e^{2G}$$

$$r(2\pi_3) = r(0) e^{3G}$$

$$\vdots$$
$$r(2\pi_n) = r(0) e^{nG}$$

$n \rightarrow +\infty$

0

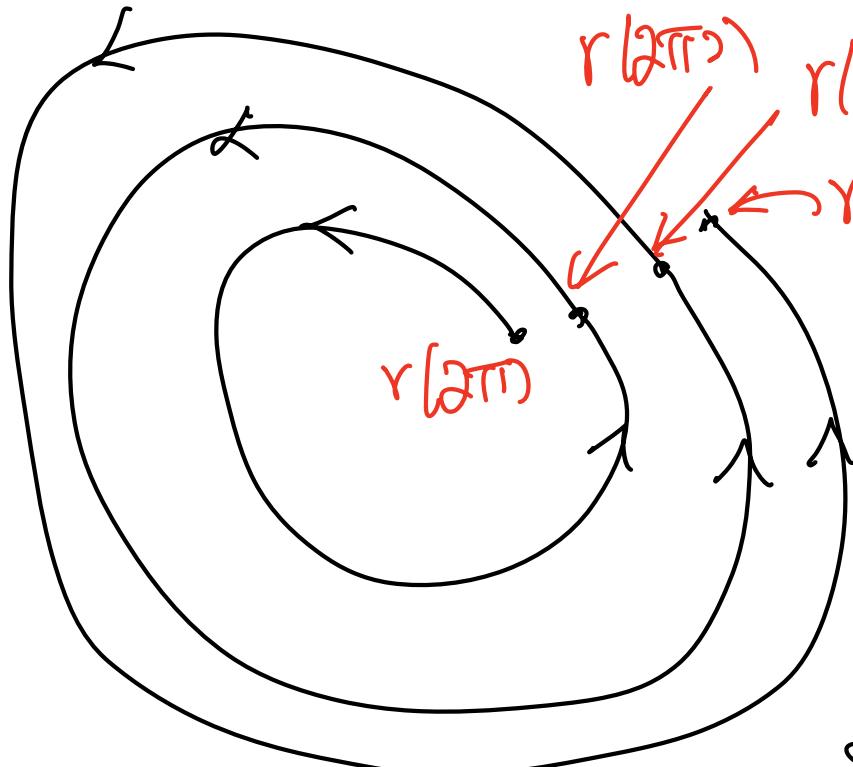


Spiral in, asymptotically stable

Solution in Polar Coordinate

Suppose $G = \int_0^{2\pi} g(\phi) d\phi$ is well-defined.

If $G > 0$, then $\ln r(2\pi) = \ln r(0) + G$



$$r(2\pi) = r(0) e^G$$

$$r(2\pi_2) = r(0) e^{2G}$$

$$r(2\pi_3) = r(0) e^{3G}$$

$$\vdots$$
$$r(2\pi_n) = r(0) e^{nG}$$

$n \rightarrow +\infty$
 $\rightarrow +\infty$

Spiral out, asymptotically unstable

Solution in Polar Coordinate

One thing to be careful

$$\left\{ \begin{array}{l} > 0 \\ < 0 \end{array} \right. ?$$

$$\frac{d\theta}{dt} = r^{n-1} [cQ(c, s) - sP(c, s)]$$

• As t increases, does $\theta(t)$ increase/decrease?

• Suppose $r(\theta) \rightarrow \begin{cases} +\infty \\ 0 \end{cases}$ as $\theta \rightarrow +\infty$

then $r(\theta(t)) \rightarrow \begin{cases} +\infty \\ 0 \end{cases}$ as $t \rightarrow +\infty$ or $-\infty$?

Solution in Polar Coordinate

Suppose $G = \int_0^{2\pi} g(\phi) d\phi$ does not exist.

How can G fail to exist?

$$g(\theta) = \frac{\cos P(\cos \theta, \sin \theta) + \sin Q(\cos \theta, \sin \theta)}{\sin Q(\cos \theta, \sin \theta) - \cos P(\cos \theta, \sin \theta)}$$



denominator becomes zero at some pt θ_c

$$\exists \theta_c \text{ s.t. } C(\theta_c)Q(\theta_c) - S(\theta_c)P(\theta_c) = 0$$

Solution in Polar Coordinate

Suppose $G = \int_0^{2\pi} g(\varphi) d\varphi$ does not exist.

How can G fail to exist?

$$g(\theta) = \frac{\cos P(\cos \theta, \sin \theta) + \sin Q(\cos \theta, \sin \theta)}{\sin Q(\cos \theta, \sin \theta) - \cos P(\cos \theta, \sin \theta)}$$

$$\int_{\theta_0}^{\theta_c} g(\varphi) d\varphi = \begin{cases} +\infty & \Rightarrow \ln r(\theta_c) = +\infty \\ & \underline{r(\theta_c) = +\infty} \\ -\infty & \Rightarrow \ln r(\theta_c) = -\infty \\ & \underline{r(\theta_c) = 0} \end{cases}$$

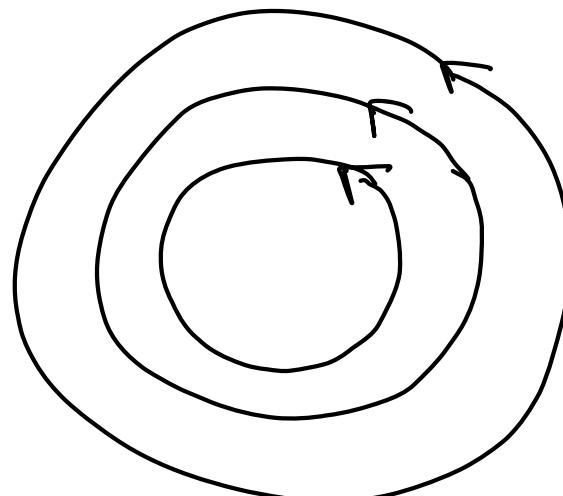
Example 1M, p.192, Ex 6.2)

$$\dot{x} = P(x, y) = -x^2y - y^3$$

$$\dot{y} = Q(x, y) = x^3 + xy^2$$

$$\dot{r} = r^3 [c(-c^2s - s^3) + s(c^3 + cs^2)] = 0$$

$$\dot{\theta} = \frac{r^3}{r} [c(c^3 + cs^2) - s(-c^2s - s^3)] = r^2 > 0$$



$$\dot{\theta} = r^2 > 0$$

Periodic Orbit

homogeneous of degree 3

periodic orbit

angle rotates

Example (M, p.192, Ex 6.3)

$$\dot{x} = P(x, y) = -(x^2 + y^2)(x+y)$$

homog. deg 2

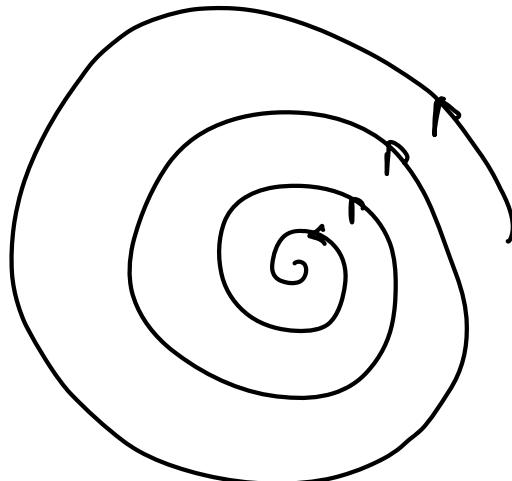
$$\dot{y} = Q(x, y) = (x^2 + y^2)(x-y)$$

$r \downarrow 0$

$$\dot{r} = r^3 \left[-c(c^2 + s^2)(c+s) + s(c^2 + s^2)(c-s) \right] = -r^3 < 0$$

$$\dot{\theta} = \frac{r^3}{r} \left[c(c^2 + s^2)(c-s) + s(c^2 + s^2)(c+s) \right] = r^2 > 0$$

θ rotate



Spiral in, asymptotically stable.

Example (M, p. 193, Ex 6.4)

$$\left\{ \begin{array}{l} \dot{x} = y^2 - x^2 \\ \dot{y} = -2xy \end{array} \right. \quad \text{homog. deg} = 2$$

$$\left\{ \begin{array}{l} \dot{r} = c[s^2 - c^2] + s[-2cs] = -r^2 \cos \theta \\ \dot{\theta} = r[c(-2cs) - s(s^2 - c^2)] = -rs \sin \theta \end{array} \right.$$

$$\frac{dr}{d\theta} = r \frac{(\cos \theta)}{(\sin \theta)}$$

$\Rightarrow 0$ at $\theta = 0, \pi$, not integrable

Example (M, p. 192, Ex 6.4)

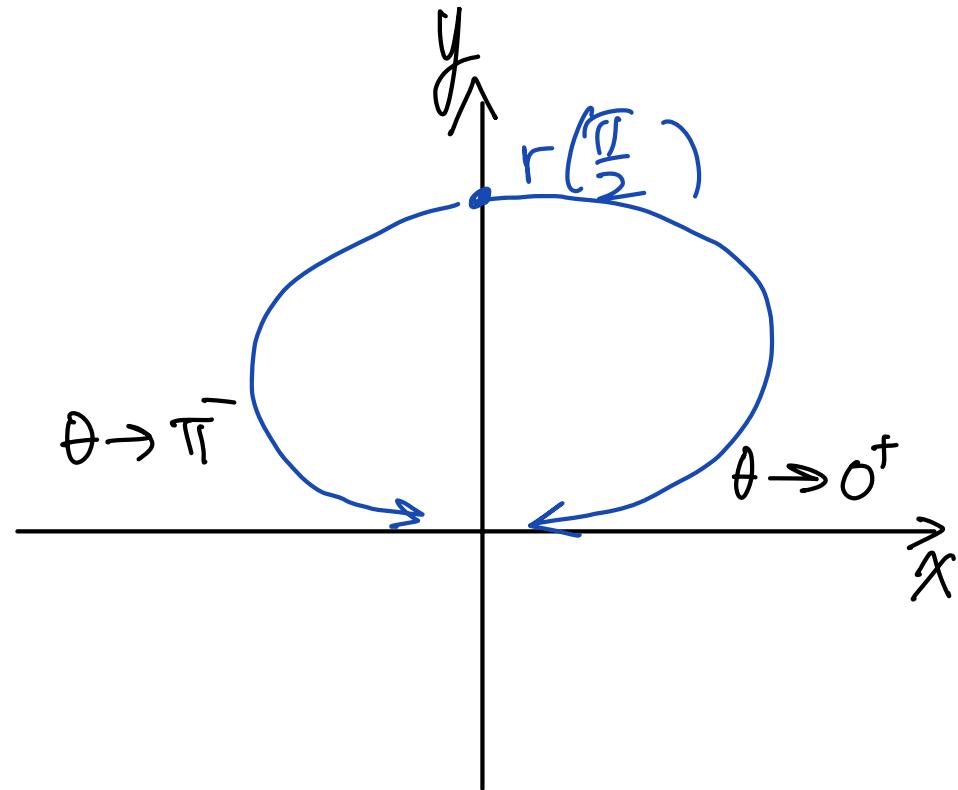
$$\ln r(\theta) = \ln r\left(\frac{\pi}{2}\right) + \int_{\frac{\pi}{2}}^{\theta} \frac{\cos \varphi}{\sin \varphi} d\varphi$$

$$\int_{\frac{\pi}{2}}^{0^+} \frac{\cos \varphi}{\sin \varphi} d\varphi = -\infty$$

$\Rightarrow r(0^+) = 0$

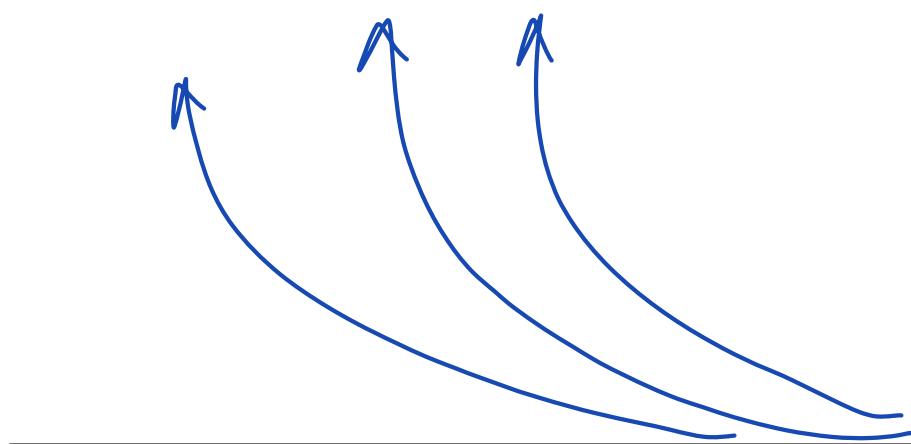
$$\int_{\frac{\pi}{2}}^{\pi^-} \frac{\cos \varphi}{\sin \varphi} d\varphi = -\infty$$

$\Rightarrow r(\pi^-) = 0$

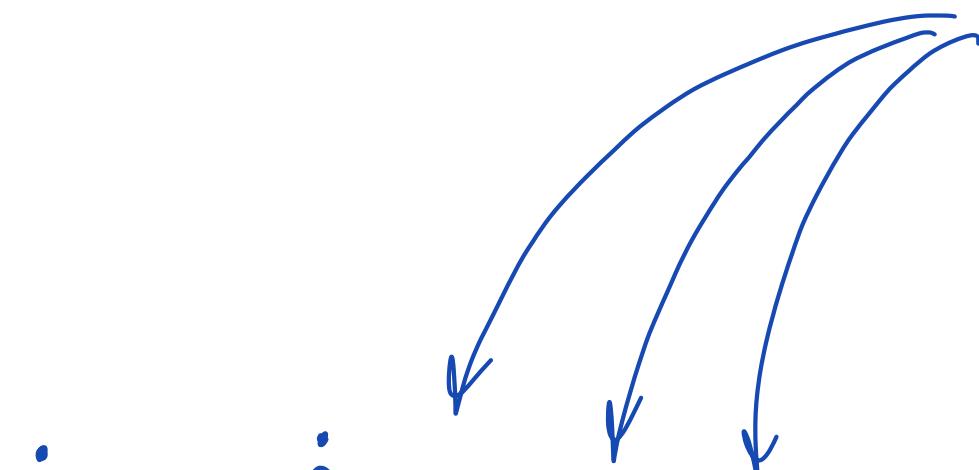


Example (M, p. 192, Ex 6.4)

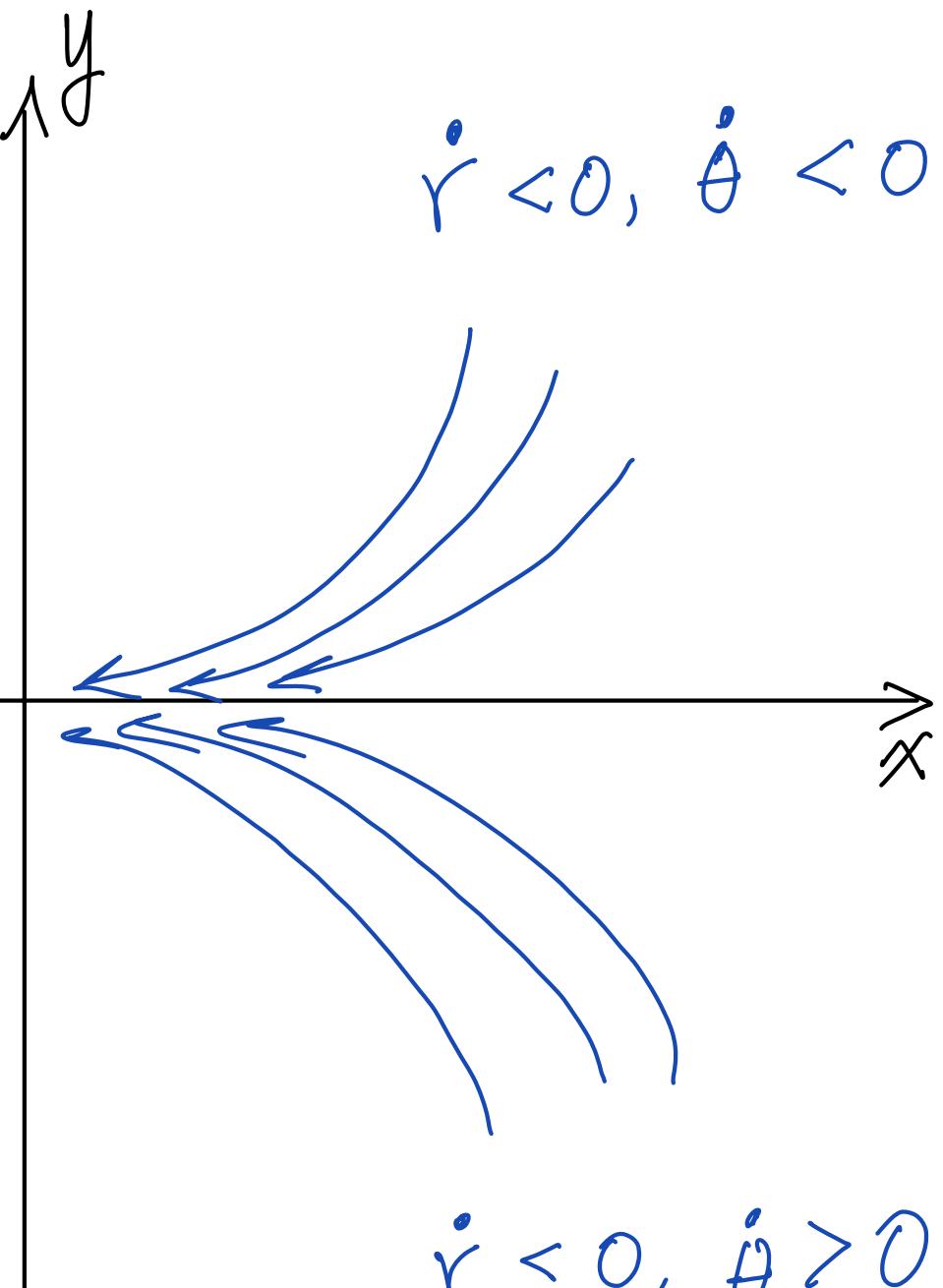
$$\dot{r} > 0, \dot{\theta} < 0$$



$$\dot{r} > 0, \dot{\theta} > 0$$



$$\dot{r} < 0, \dot{\theta} < 0$$



$$\dot{r} < 0, \dot{\theta} > 0$$

Example (M, p. 192, Ex 6.4)

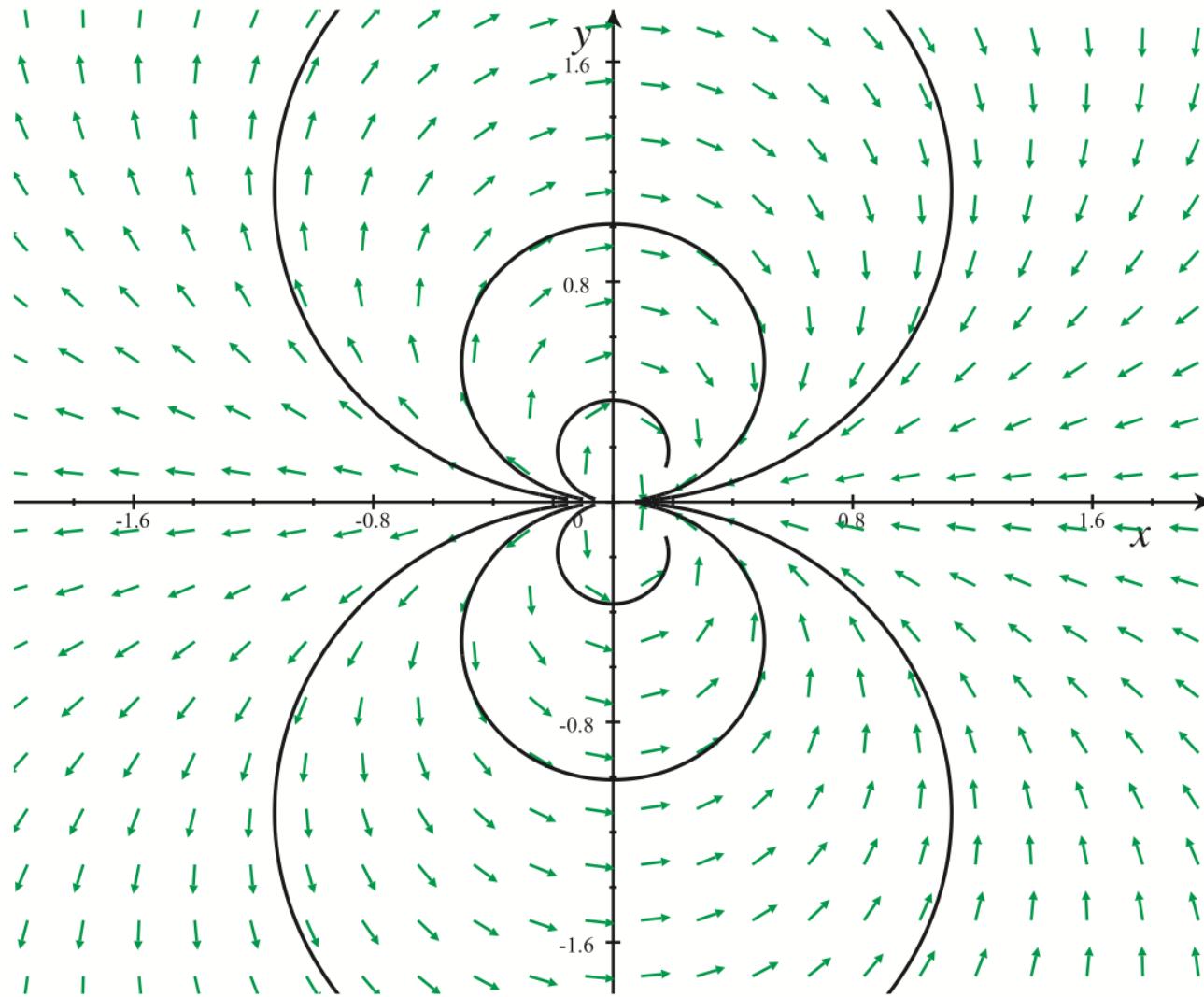


Figure 6.1. Phase portrait of the example (6.3).

Example (M, p.194, Ex 6.5)

$$\begin{cases} \dot{x} = y^2x - x^2y \\ \dot{y} = x^3 + y^3 \end{cases}$$

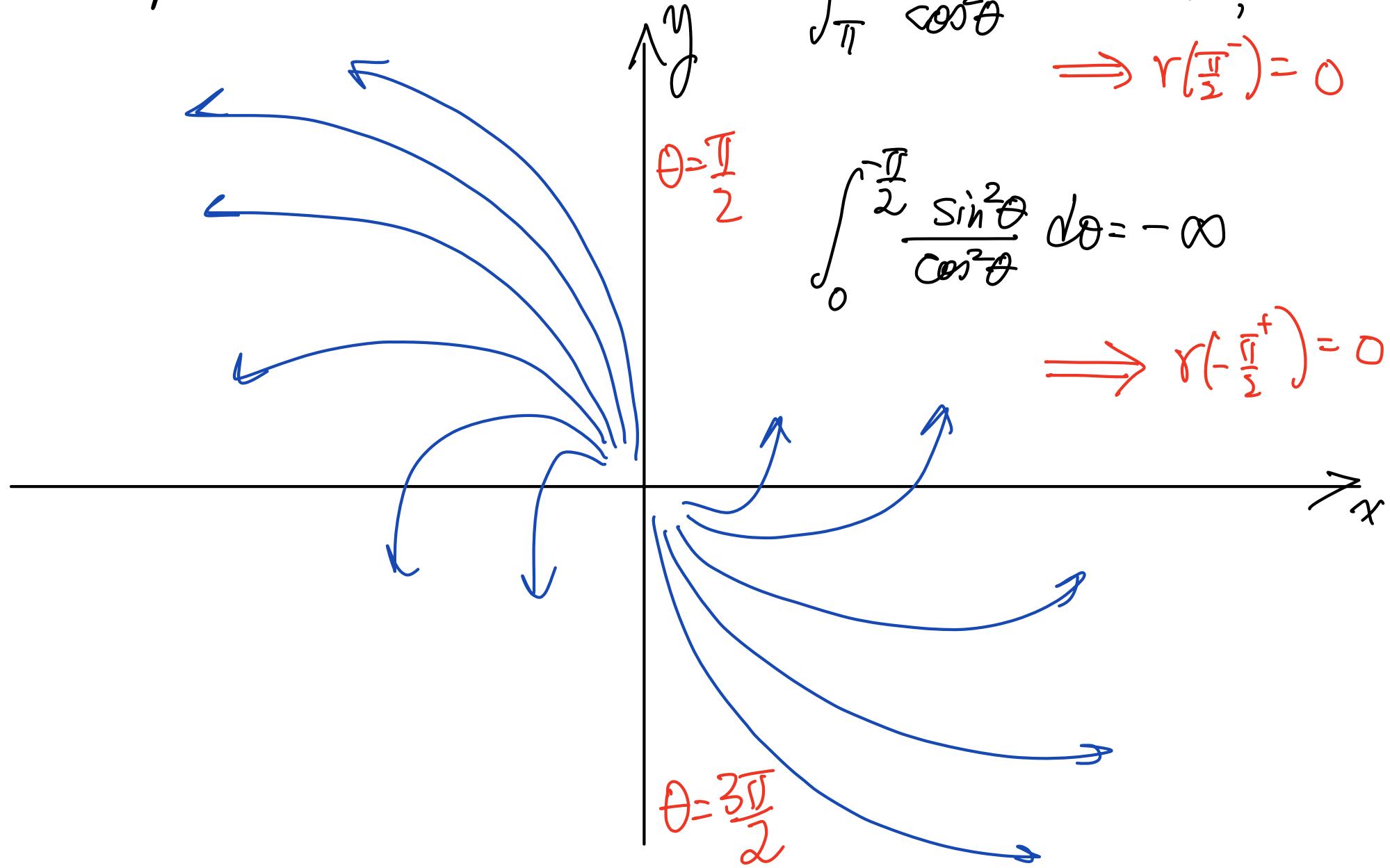
$$\dot{r} = r^3 [c(s^2c - c^2s) + s(c^3 + s^3)] = r^3 \sin^2\theta > 0$$

$$\dot{\theta} = \frac{r^3}{r} [c(c^3 + s^3) - s(s^2c - c^2s)] = r^2 \cos^2\theta > 0$$

$$\frac{dr}{d\theta} = r \frac{\sin^2\theta}{\cos^2\theta}$$

$\cos^2\theta = 0$ at $\theta = \frac{\pi}{2}, \frac{3\pi}{2}$

Example (M, p. 194, Ex 6.5)



Example (M, p.194, Ex 6.5)

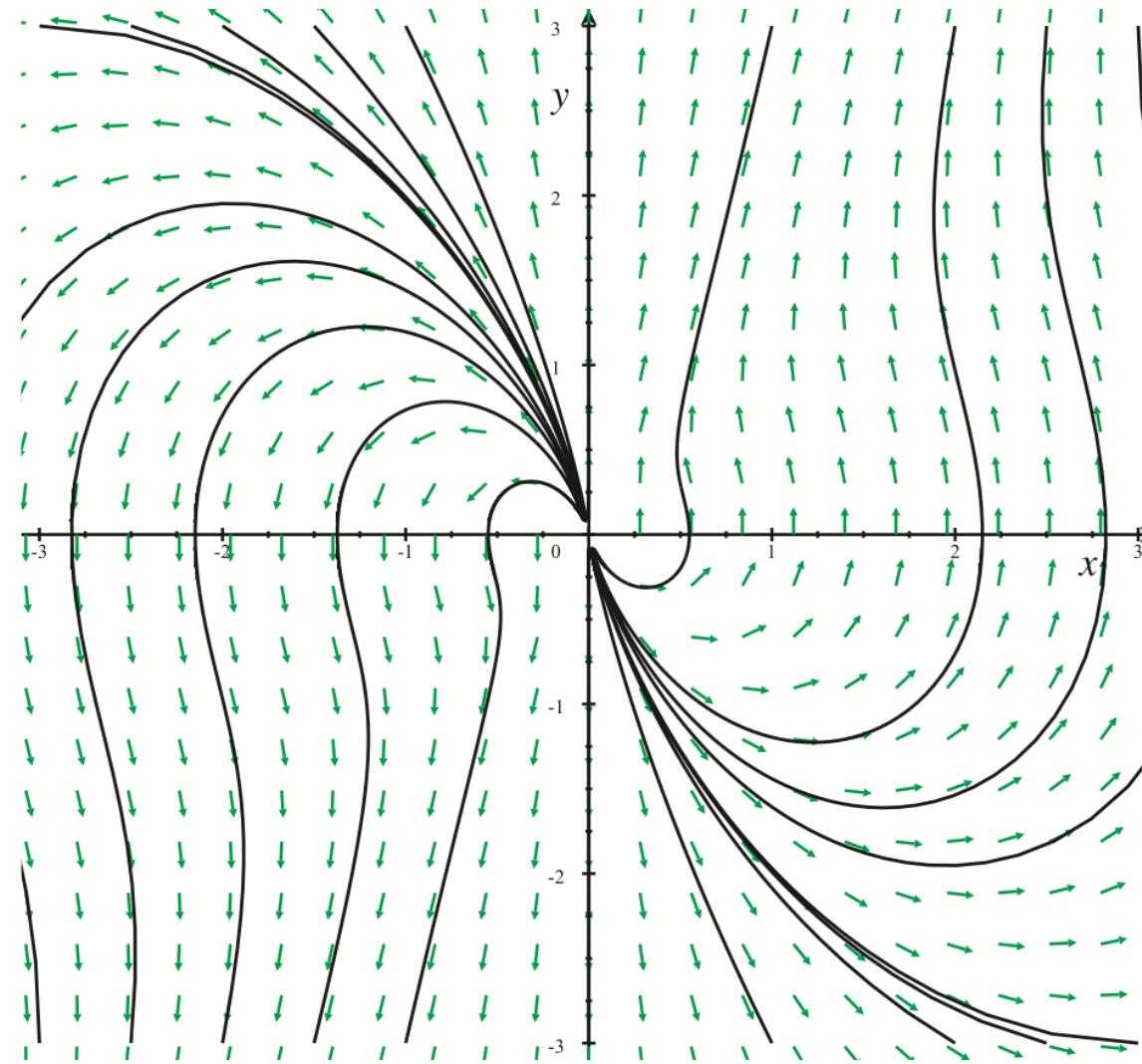


Figure 6.3. Phase portrait of (6.11).

Complete characterization (Case ③)

(M P. 192) Near o (inside $B_\delta(o)$ for some δ)

- ① Topological center: every orbit is periodic enclosing o .
- ② Nonhyperbolic focus: $r(t) \rightarrow o$, $\theta(t) \rightarrow +\infty$ as either $t \rightarrow +\infty$ or $t \rightarrow -\infty$
- ③ Nonhyperbolic node: \exists some θ_c s.t. $g(\theta_c) = 0$
 $r(\theta) \rightarrow +\infty$ or $-\infty$ as $\theta \rightarrow \theta_c^\pm$

Case ④, linear center $A = \begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix}$

$$\begin{cases} \dot{x} = -\omega y + p(x, y) \\ \dot{y} = \omega x + q(x, y), \quad |p, q| \leq C(x^2 + y^2) \end{cases}$$

$$\begin{aligned} \dot{r} &= c \cancel{[-\omega r s + p(r_c, r_s)]} + s \cancel{[\omega r c + q(r_c, r_s)]} \\ &= c p(r_c, r_s) + s q(r_c, r_s) \leq \underline{\mathcal{O}(r^2)} \end{aligned}$$

$$\begin{aligned} \dot{\theta} &= \frac{1}{r} \left[c \cancel{(\omega r c + q(r_c, r_s))} - s \cancel{(-\omega r s + p(r_c, r_s))} \right] \\ &= \omega + \frac{1}{r} [c q(r_c, r_s) - s p(r_c, r_s)] = \underline{\omega + \mathcal{O}(r)} \end{aligned}$$

Case ④, linear center $A = \begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix}$

$$\begin{cases} \dot{x} = -\omega y + p(x, y) \\ \dot{y} = \omega x + q(x, y), \quad |p, q| \leq C(x^2 + y^2) \end{cases}$$

$$\dot{r} = c p(r_c, r_s) + s q(r_c, r_s) \leq \underline{\partial(r^2)}$$

$$\dot{\theta} = \omega + \frac{1}{r} [c q(r_c, r_s) - s p(r_c, r_s)] = \underline{\omega + O(r)}$$

Note: for $r \ll 1$, $\dot{\theta} \approx \omega$, i.e. θ rotates with speed ω

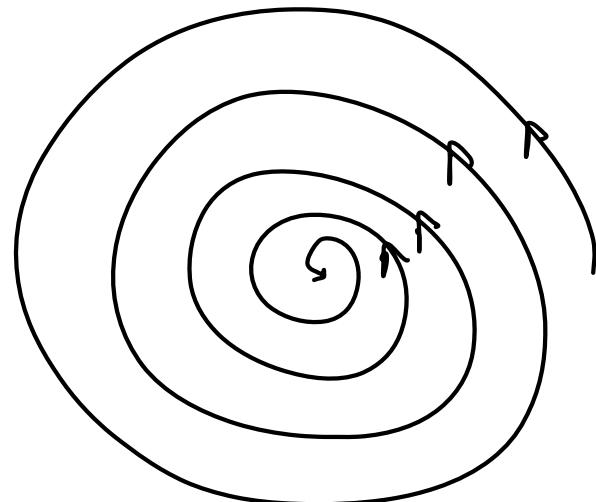
Example (p, 197, Ex 6.8)

$$\begin{cases} \dot{x} = -\omega y - x^3 \\ \dot{y} = \omega x - yx^2 \end{cases}$$

$$g(x, y) = -x^2$$

$$\begin{pmatrix} \dot{x} = -\omega y + x g(x, y) \\ \dot{y} = \omega x + y g(x, y) \end{pmatrix}$$

$$\begin{cases} \dot{r} = c(-r^3 c^3) + s(-r^3 s c^2) = -r^3 c \cos^2 \theta < 0 \\ \dot{\theta} = \omega + \frac{r^3}{r} [c(-s c^2) - s(c^3)] = \omega \end{cases}$$



spiral in,
asymptotically stable

Example (p. 198 Ex 6.10)

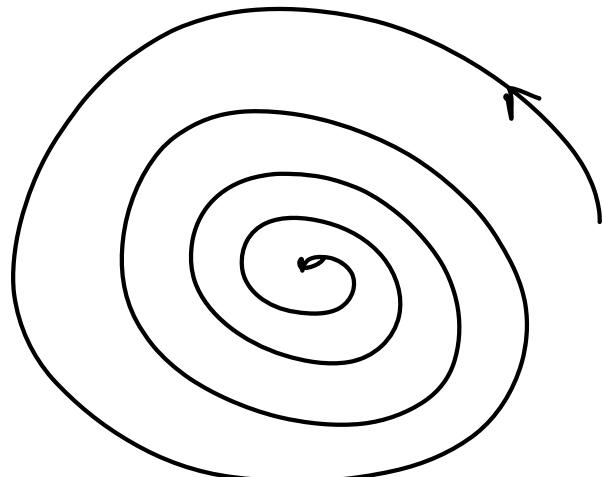
$$\begin{cases} \dot{x} = -y - x^3 - xy^2 \\ \dot{y} = x - yx^2 - y^3 \end{cases}$$

$$g(x,y) = -x^2 - y^2$$

$$\begin{cases} \dot{x} = -\omega y + x g(x,y) \\ \dot{y} = \omega x + y g(x,y) \end{cases}$$

$$\dot{r} = r^3 [c(-c^3 - cs^2) + s(-sc^2 - s^3)] = -r^3 < 0$$

$$\dot{\theta} = 1 + r^2 [c(-sc^2 - s^3) - s(-c^3 - cs^2)] = 1$$



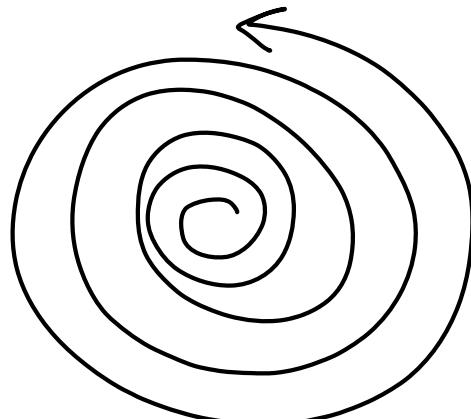
Spiral in,
asymptotically stable

Example (p. 199, Ex 6.11)

$$\begin{cases} \dot{x} = -\omega y + xy^2 + \alpha^2 y + y^3 \\ \dot{y} = \omega x + y^3 - x^3 - xy^2 \end{cases}$$

$$\dot{r} = r^3 [c(c s^2 + c^2 s + s^3) + s(s^3 - c^3 - c s^2)] = r^3 \sin^2 \theta \geq 0$$

$$\dot{\theta} = \omega + r^2 [c(s^3 - c^3 - c s^2) - s(c s^2 + c^2 s + s^3)] = \underbrace{\omega - r^2}_{> 0} \quad \text{for } r \ll 1$$



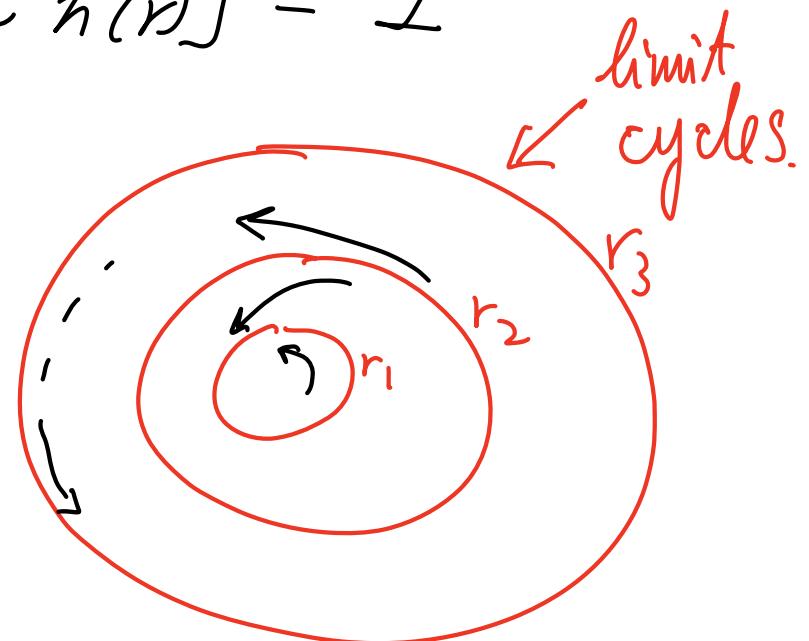
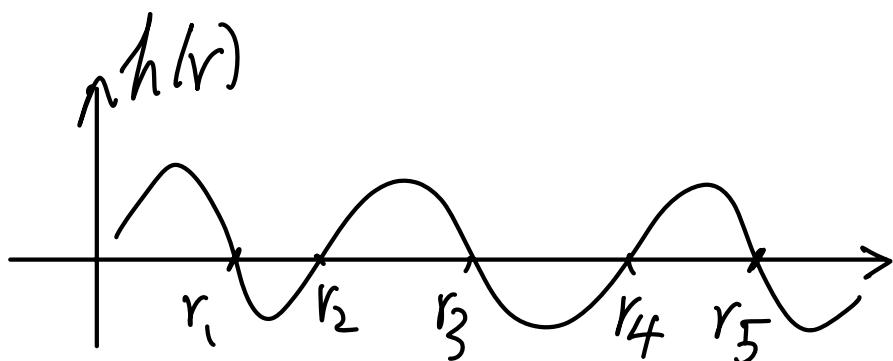
spiral out,
unstable

Example (P. 201, Ex 6.14)

$$\begin{cases} \dot{x} = -y + xh(r) \\ \dot{y} = x + yh(r) \end{cases} \quad (r = \sqrt{x^2 + y^2})$$

$$\dot{r} = c[r]ch(r) + s[r]sh(r) = rh(r)$$

$$\dot{\theta} = 1 + c[s]h(r) - s[c]h(r) = 1$$



Complete characterization of linear center

(4)

(M P. 197) Near o (inside $B_\delta(o)$ for some δ)

- ① Topological center: every orbit is periodic enclosing o .
- ② Nonhyperbolic focus: $r(t) \rightarrow 0$, $\theta(t) \rightarrow +\infty$ as either $t \rightarrow +\infty$ or $t \rightarrow -\infty$
- ③ Center-focus: there is an infinite nested limit cycles $\gamma_n \xrightarrow{n \rightarrow \infty} o$ and every orbit between γ_n & γ_{n+1} cycles toward γ_n, γ_{n+1} as $t \rightarrow \pm\infty$